

# **A method to estimate clear sky mesoscale vertical motion from geostationary satellite imagery**

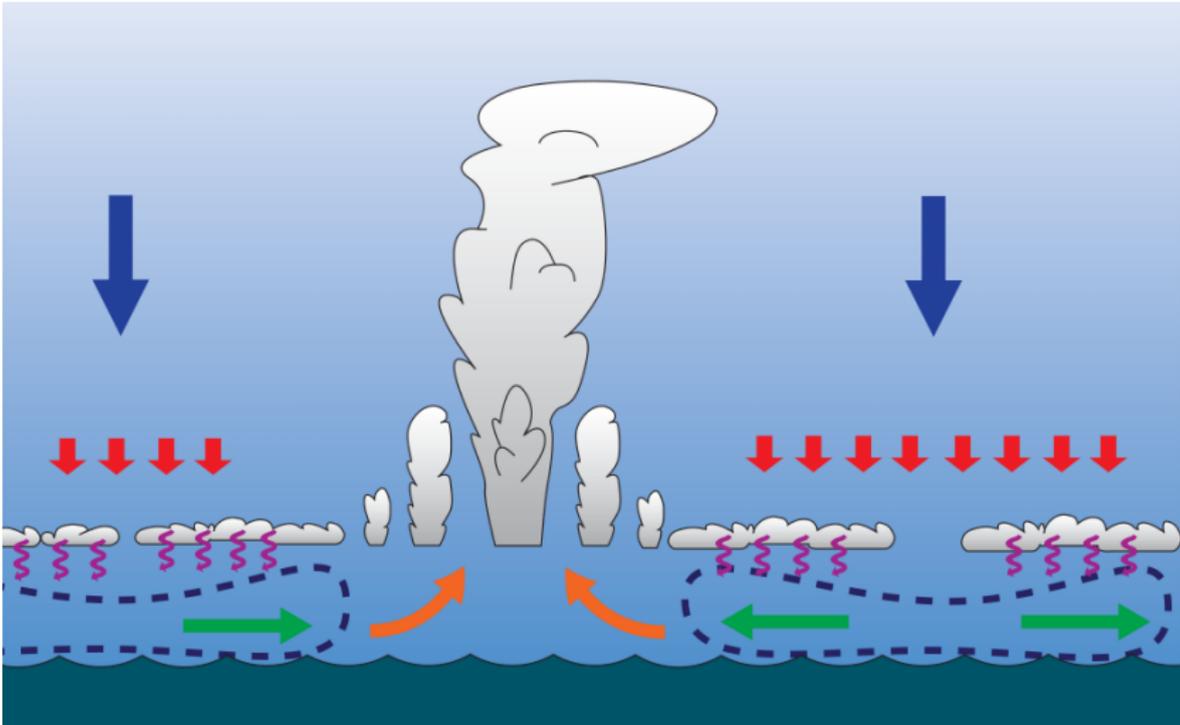
Basile Poujol and Sandrine Bony

LMD Paris

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# Motivations

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*Coppin and Bony, 2015*

- Importance of clear sky vertical motion in cloud organization and radiation-circulation coupling
- Sparsity and cost of in situ observations
- Currently no available satellite measurement
- Future satellite missions focus on assessing  $w$  within clouds

## Outline

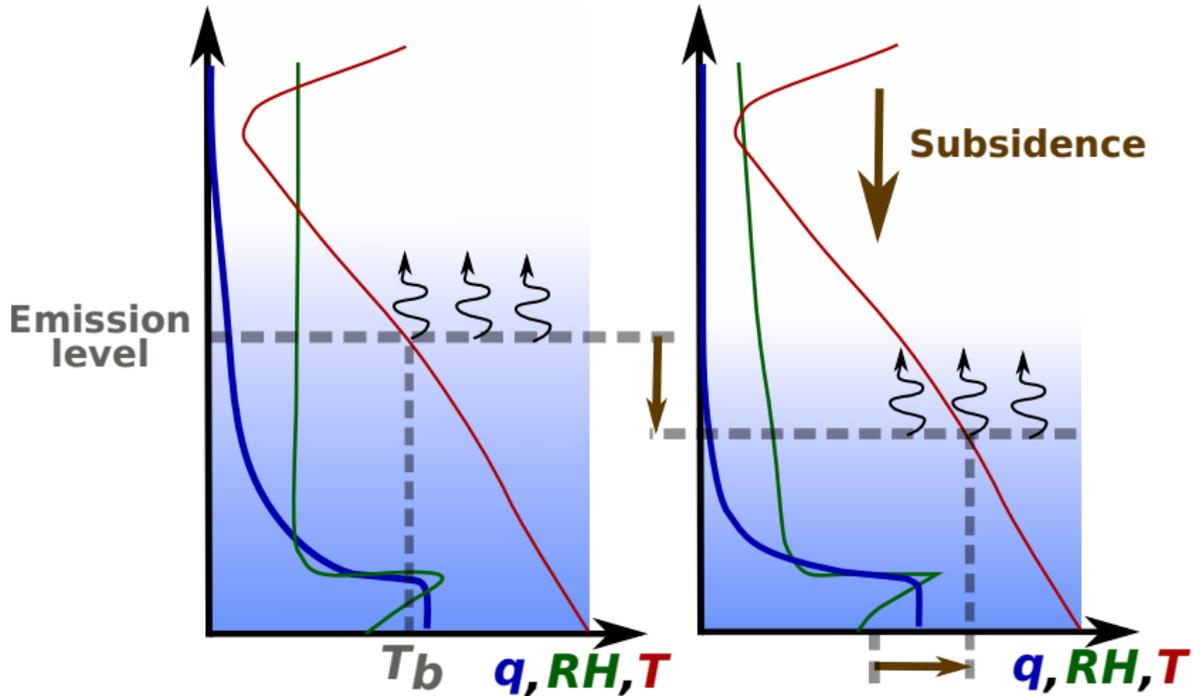
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- Theory
- Implementation
- Evaluation
- Insights on possible causes of vertical motion



# Main idea of the method

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Subsidence dries the atmosphere and increases brilliance temperature

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## Hypotheses

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- (H1) At the wavelength considered **water vapor is the only significant absorber** and  $e^{-\tau_s} \ll 1$
- (H2) **Relative humidity is vertically uniform in the vicinity of the emission level**
- (H3) The specific extinction coefficient  $\kappa$  is vertically uniform in the vicinity of the emission level
- (H4) Moist adiabatic lapse rate (WTG)
- (H5) **Hydrostatic approximation**
- (H6) Perfect gas
- (H7) No water vapour turbulent fluxes

## Link between $RH$ and temperature at emission level

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Formula for optical thickness :

$$\tau = \int_{TOA}^p \frac{\kappa}{g} q dp$$

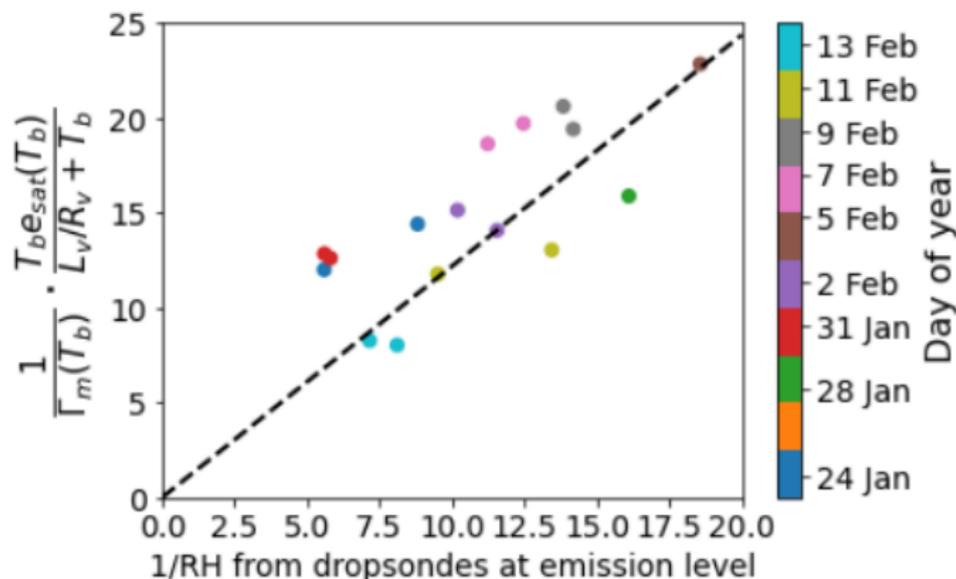
Integration done using :

- hydrostatic approximation
- perfect gas law
- Clausius Clapeyron law

## Link between $RH$ and temperature $T^*$ at emission level

Result :

$$\frac{R_d T}{\epsilon \kappa} \times \frac{1}{RH} = \frac{1}{\Gamma_m(T^*)} \frac{T^* e_{sat}(T^*)}{L_v/R_v + T^*}$$



## Link between relative humidity variations and vertical motion

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Conservation of specific humidity :

$$q(p, t + dt) = q(p - \omega dt, t)$$

which implies a direct link between vertical velocity and relative humidity variations.

$$\frac{dRH}{dt} = -\frac{RH}{p} \left( \frac{L_s R_d \Gamma_m}{R_v g T} - 1 \right) \times \omega$$

**Direct link between vertical velocity  $\omega$  and variations of temperature  $T^*$  at emission level :**

$$\omega = \frac{1}{\frac{L_s R_d \Gamma_m}{R_v g T^*} - 1} \frac{L_s p}{R_v T^{*2}} \frac{L_s / R_v + 2T^*}{L_s / R_v + T^*} \times \frac{dT^*}{dt}$$

The relation is independent on spectroscopic properties !

## Implementation main steps

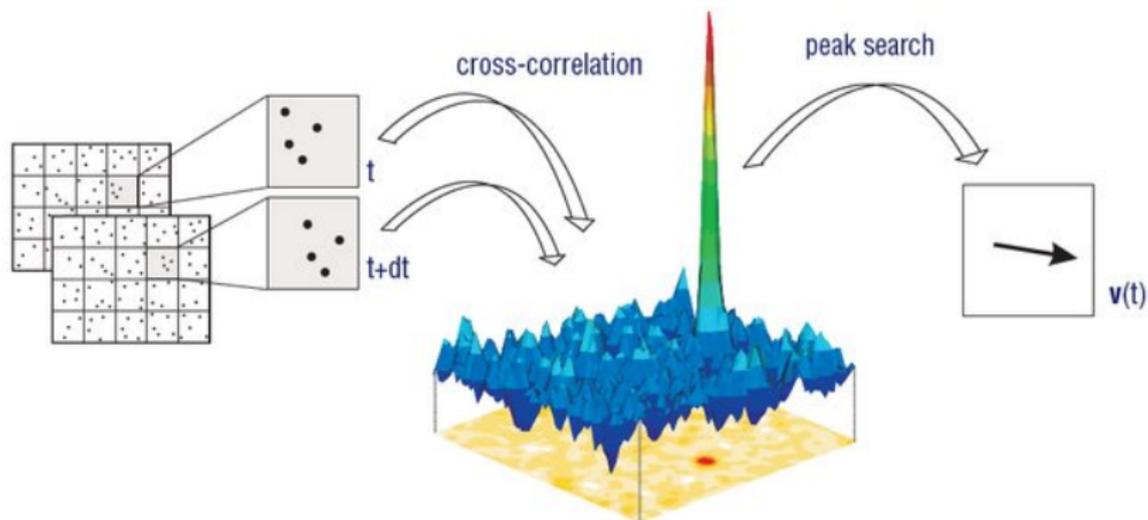
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- Compute horizontal winds
- Compute temperature  $T^*$  from satellite radiances
- Compute Lagrangian derivative  $dT^*/dt$
- Compute vertical velocity

## Determination of horizontal winds

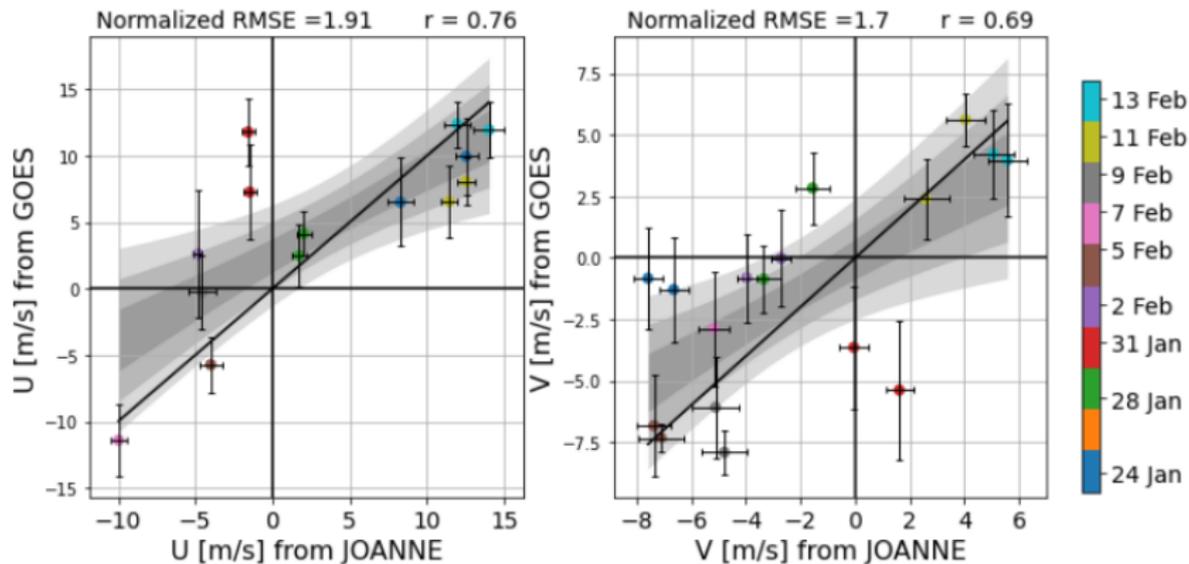
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Per-filtering to select water vapour filaments.



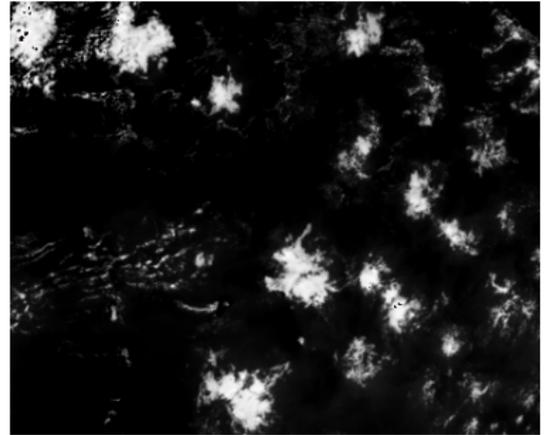
*Wieneke, 2017*

# Determination of horizontal winds



## Determination of temperature at $\tau = \tau^*$

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Using these images we deduce :

- Optical depth  $\tau_s$  of the atmosphere at wavelength  $\lambda$
- Sea Surface / Cloud Top Temperature  $T_{sfc}$

Radiative transfer equation :

$$R_\lambda = \int_0^{\tau_s} B_\lambda(\tau) e^{-\tau} d\tau + e^{-\tau_{sfc}} B_\lambda(T_{sfc})$$

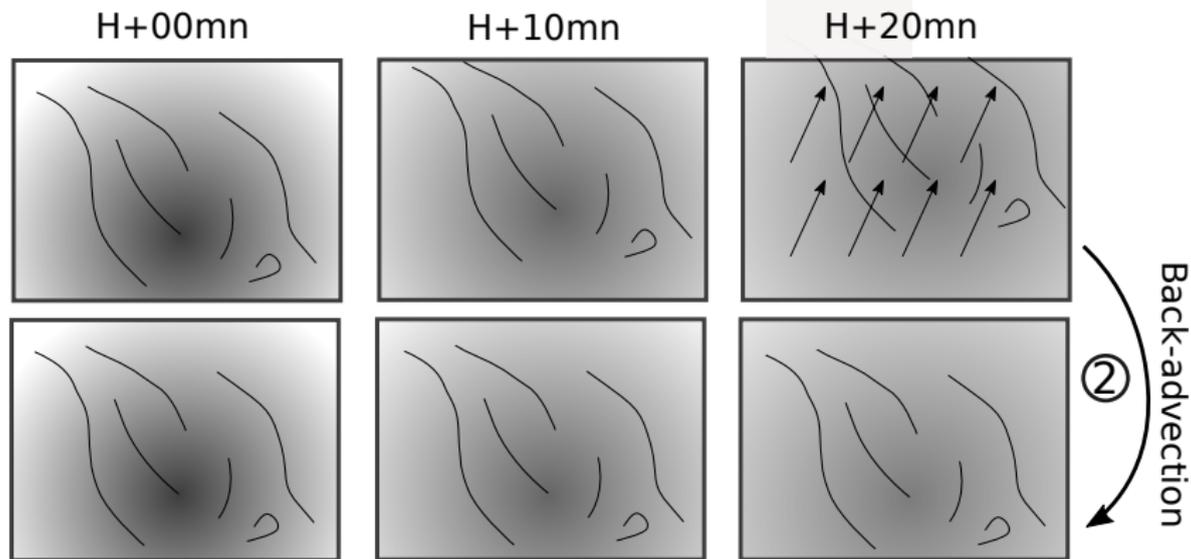
We invert this equation

→ We can deduce temperature  $T^* = f(R_\lambda, T_{sfc})$

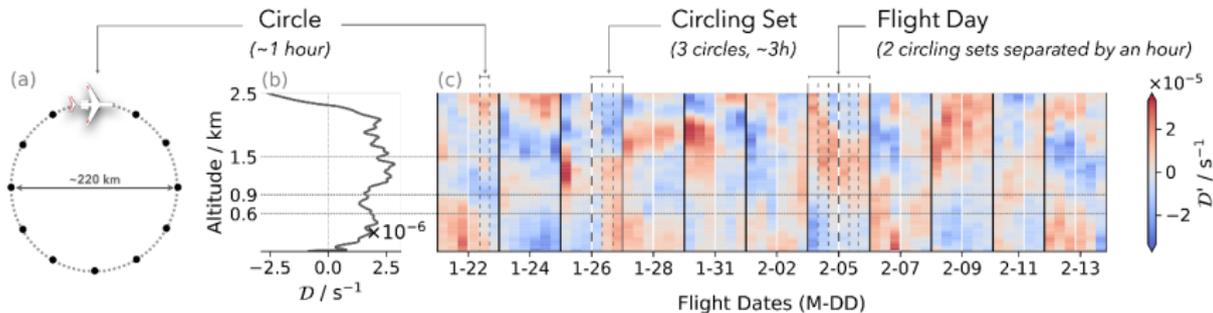
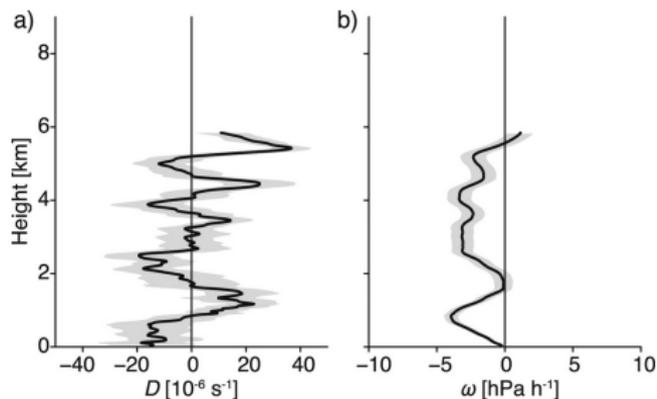
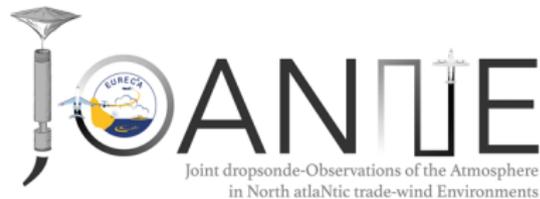
→ Correction from surface or cloud-top temperature inhomogeneities

# Determination of Lagrangian derivative of temperature $dT^*/dt$

① Compute horizontal winds



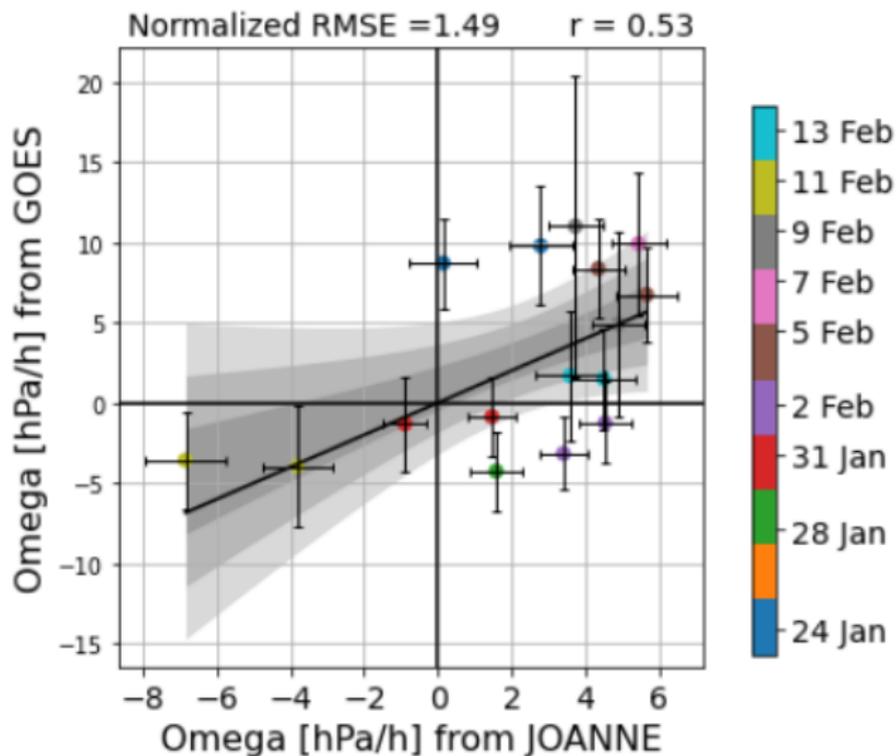
# Evaluation against JOANNE data



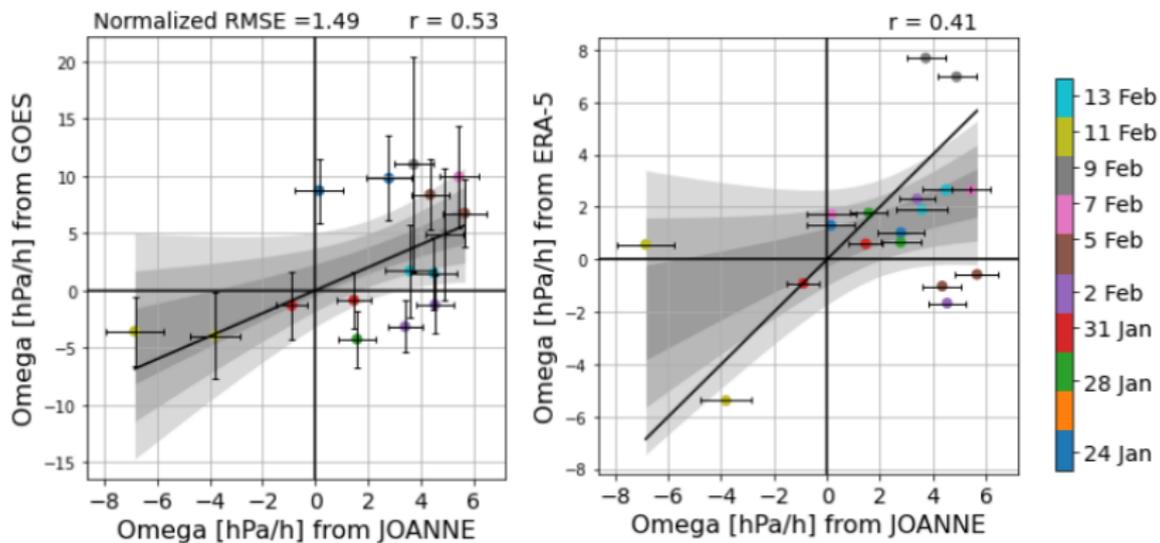
*Bony and Stevens 2019; George et al. 2021,2022*

## Evaluation against JOANNE data

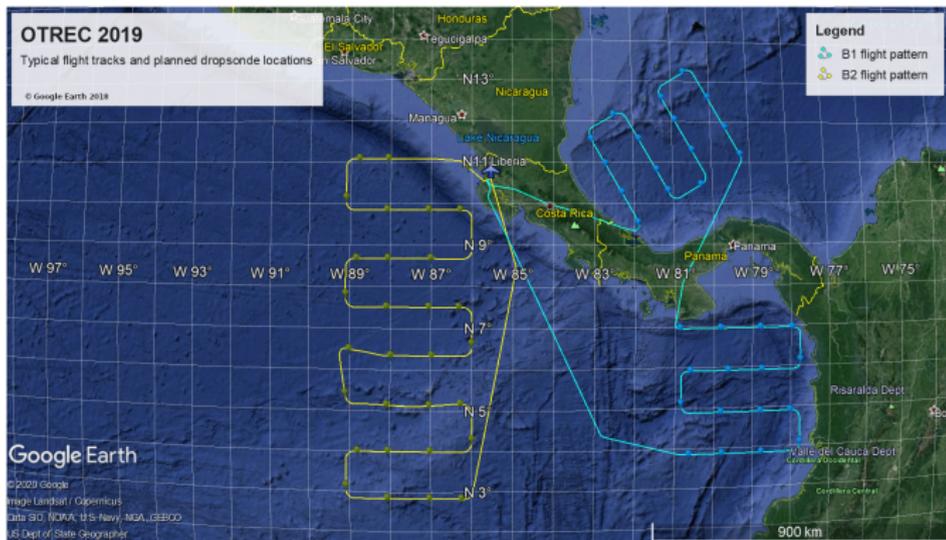
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# Evaluation against JOANNE data



## Grids of dropsondes in tropical Pacific during Aug-Oct 2019



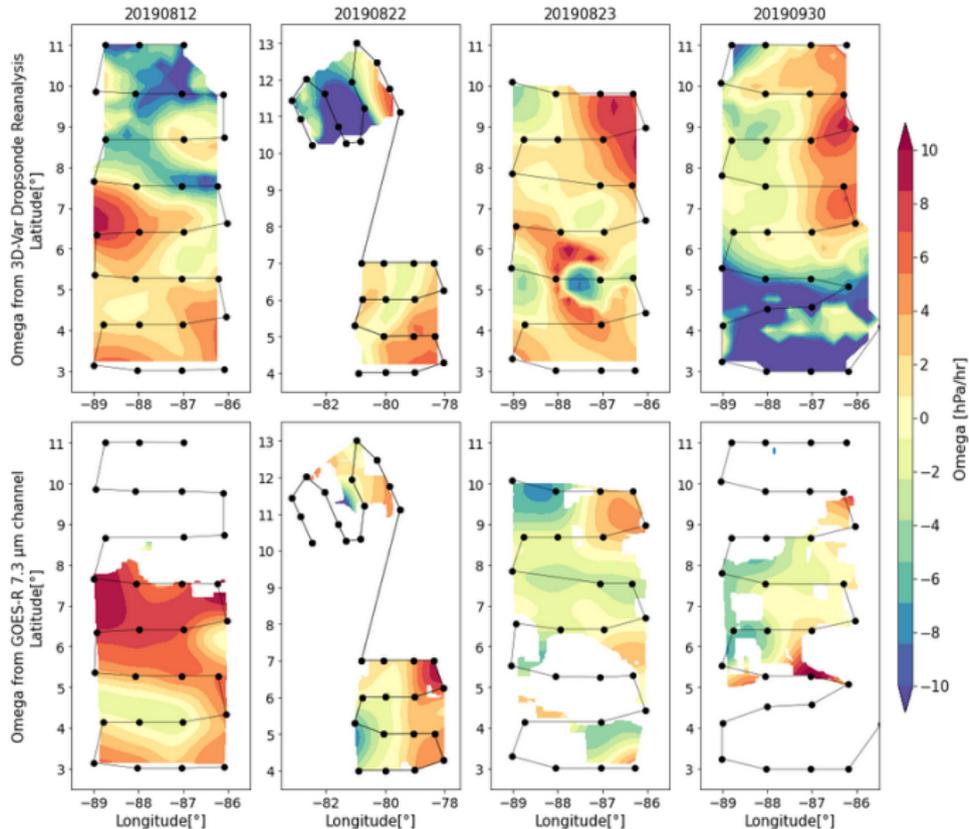
*Vömel et al., 2021*

- In regions of active deep convection
- Flight duration of 3-5h
- Dropsonde resolution 1 degree

Vertical velocity computed from :

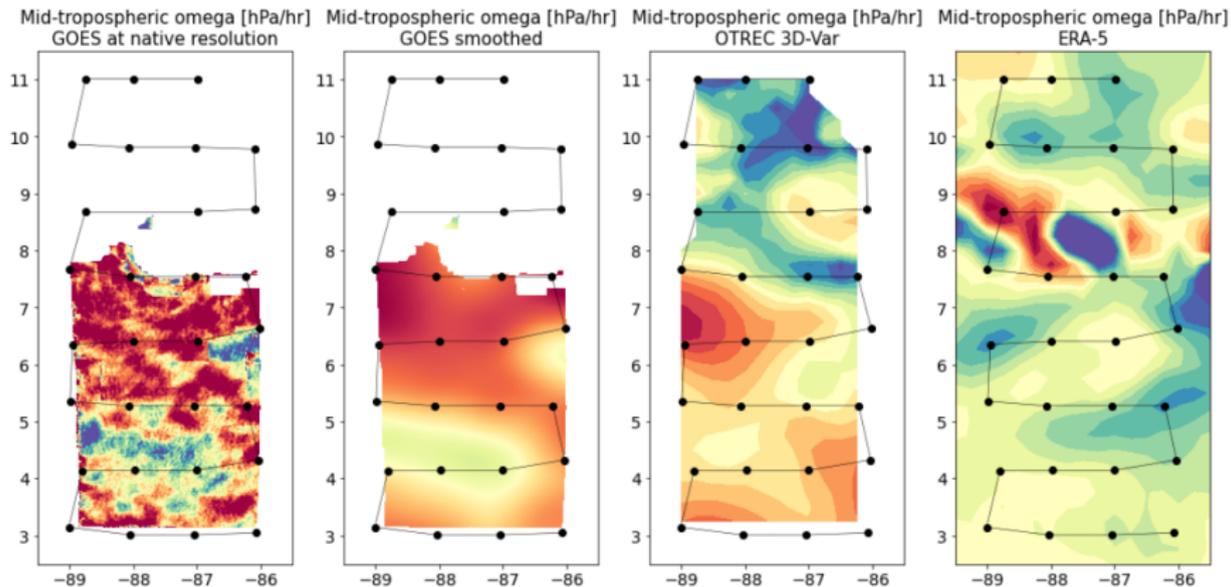
- Dropsondes using mass conservation (*Raymond & Fuchs-Stone, 2021*)
- GOES images for several hours, then interpolated onto flight track using the closest time from dropsonde launch

# Evaluation against OTREC data



# Evaluation against OTREC data

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## Evaluation in the NARVAL simulations

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- 2.5km resolution in the Atlantic ( $500 \times 2500$ km)
- NARVAL-1 : Winter trades (December 2013)
- NARVAL-2 : ITCZ edge (August 2016)
- Hourly output of 3D fields

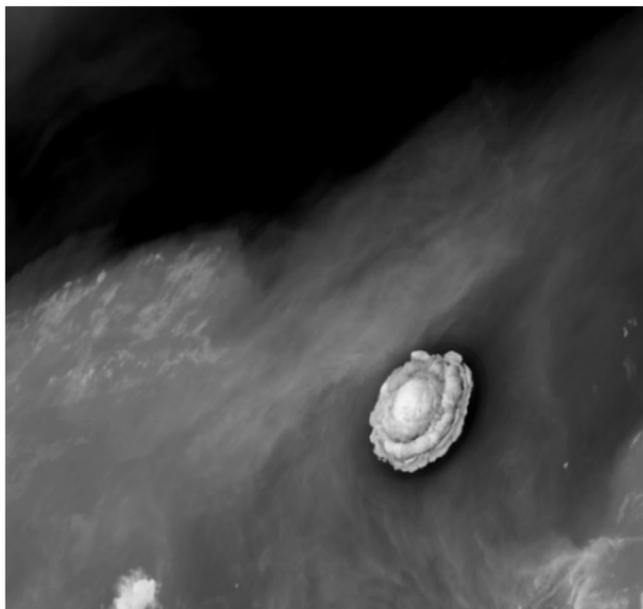
Simulation of brilliance temperature of GOES-16 ABI Instrument using radiative transfer code RTTOV.

Simulated performances are similar to those during OTREC and EUREC4A.

## Evaluation at subhourly time scale

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15 January 2022 : Hunga Tonga eruption triggered worldwide gravity waves

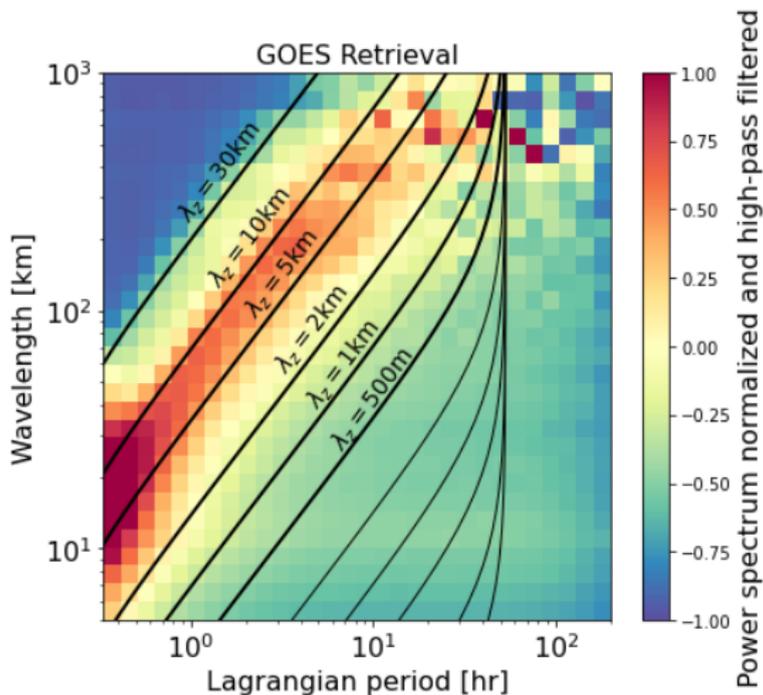


GOES-W image at  $6.9\mu\text{m}$



# Gravity waves as a good candidate for mesoscale vertical motion

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$$(\omega - \vec{u} \cdot \vec{k})^2 = N^2 \frac{\vec{k}^2}{k^2 + m^2 + \frac{1}{4H^2}} \quad (\text{inspired from Wheeler\&Kiladis 1999})$$

## Gravity waves as a good candidate for mesoscale vertical motion

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## Conclusions

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- Ability to retrieve temporal and spatial variations of mesoscale vertical velocity
- Limited quantitative performance
- Measurements at high spatial (2km) and temporal (10mn) resolution over full geostationary disk
- First results seem to support the hypothesis for gravity waves (vs radiation) as a main cause for mesoscale vertical motion

Next steps :

- Evaluation of the method against VHF radar data ?
- Use of the tool to study convective aggregation

## Determination of temperature at $\tau = \tau^*$

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**At what level  $\tau^*$  do we measure  $\frac{dT^*}{dt}$  ?**

Simple possible hypothesis :  $\tau^* = 1 \rightarrow T^* = T_b$

But actually we observe variations of temperature at the following level :

$$\tau^* = \langle \tau \rangle_{\frac{\delta T_b}{\delta T}} = 1 + \eta \approx 1.35$$

where  $\eta = \frac{hcR_v}{\lambda L_s k_B}$  depends only on the channel wavelength.

## Evaluation in the NARVAL simulations

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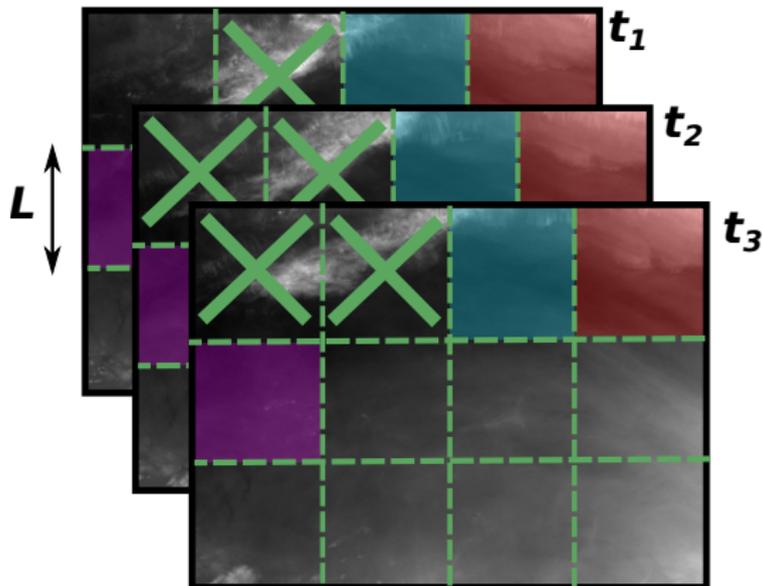
- 2.5km resolution in the Atlantic ( $500 \times 2500$ km)
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Simulation of brilliance temperature of GOES-16 ABI Instrument using radiative transfer code RTTOV.

Because output is "only" hourly it is not possible to use a linear regression to compute  $dT_b/dt$ , only a difference between H and H+1 is done.  $\rightarrow$  we expect a poorer performance of the retrieval

## Evaluation in the NARVAL simulations

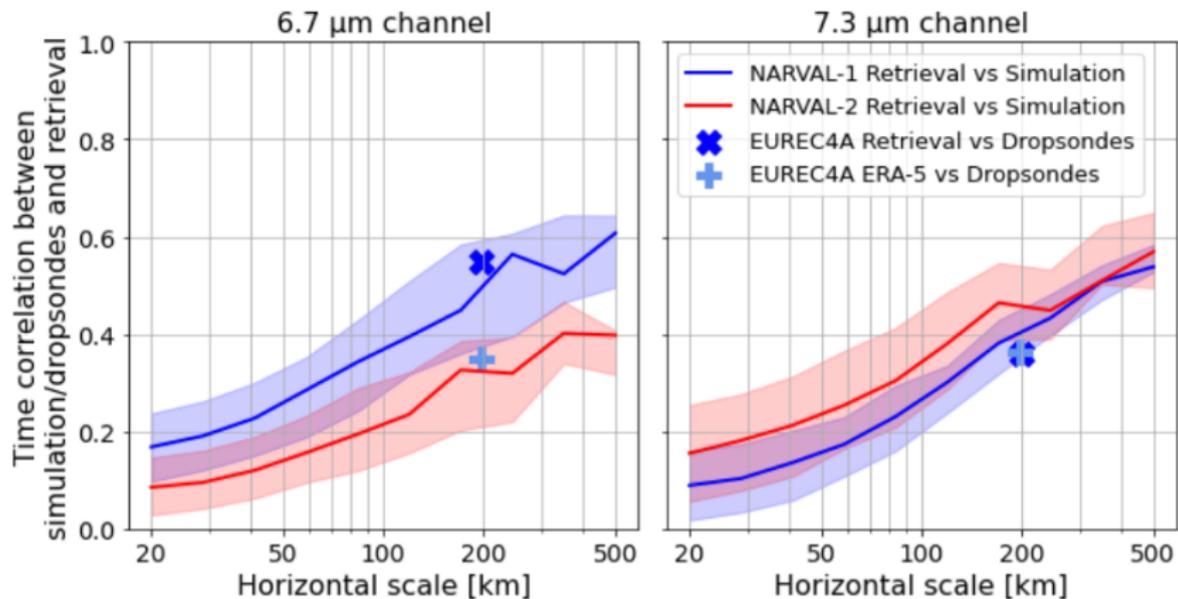
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One time series of omega in the simulation and the retrieval for each colour  $\rightarrow$  one correlation per block  $\rightarrow$  mean correlation and uncertainty

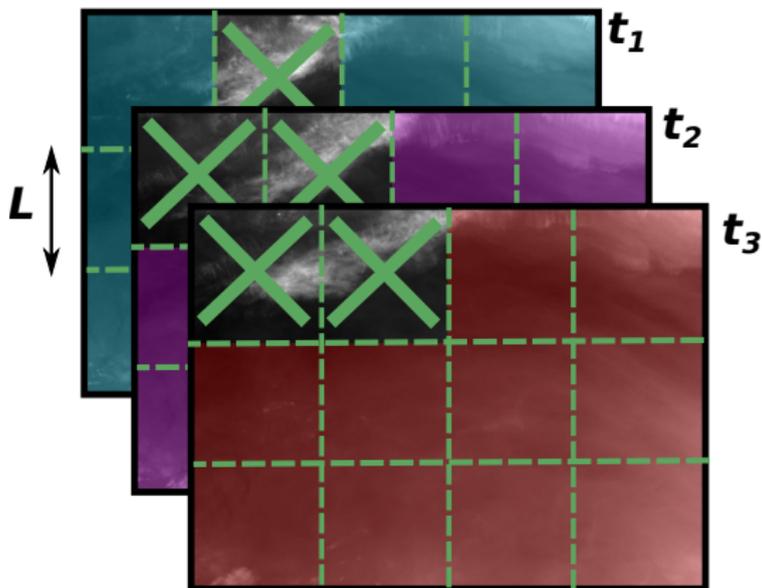
## Evaluation in the NARVAL simulations

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## Evaluation in the NARVAL simulations

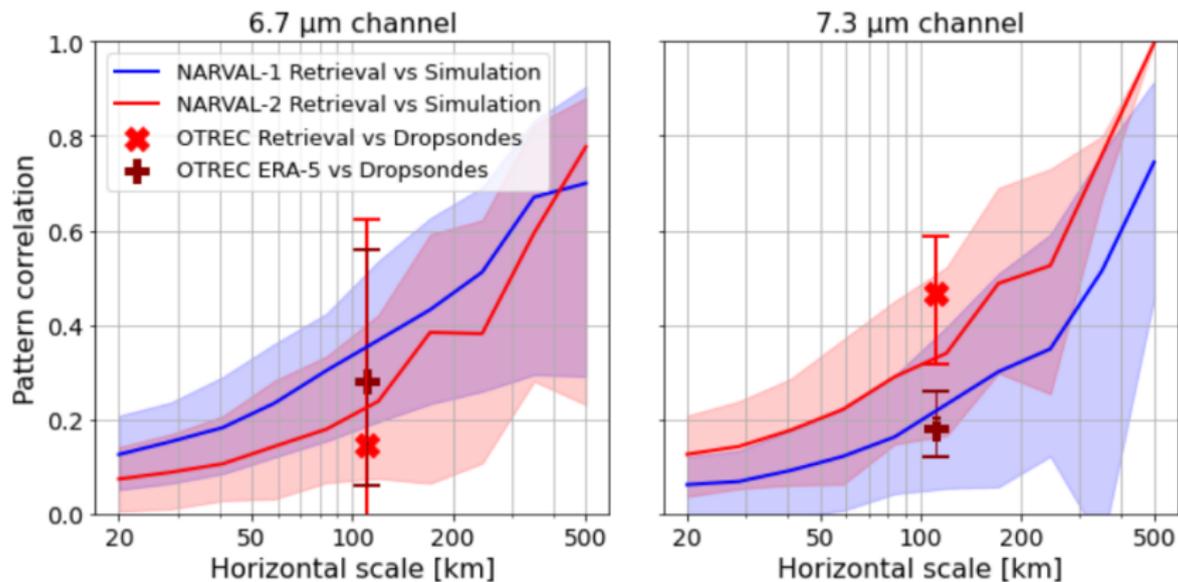
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One series of omega in the simulation and the retrieval for each colour  
→ one correlation per timestep → mean correlation and uncertainty

# Evaluation in the NARVAL simulations

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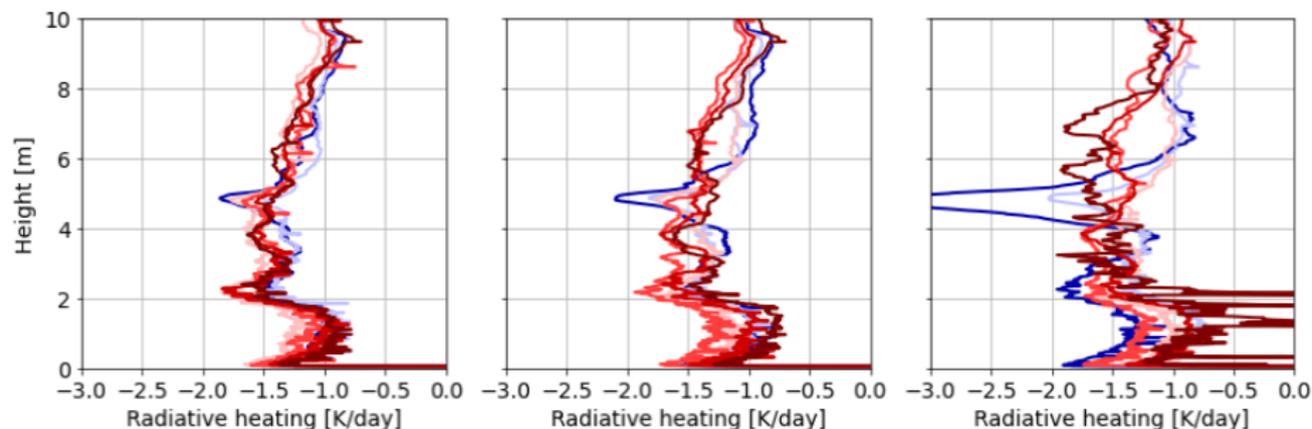
## Gravity waves as a good candidate for mesoscale vertical motion

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Small scale(10km)

Mesoscale(100km)

Large scale(500km)

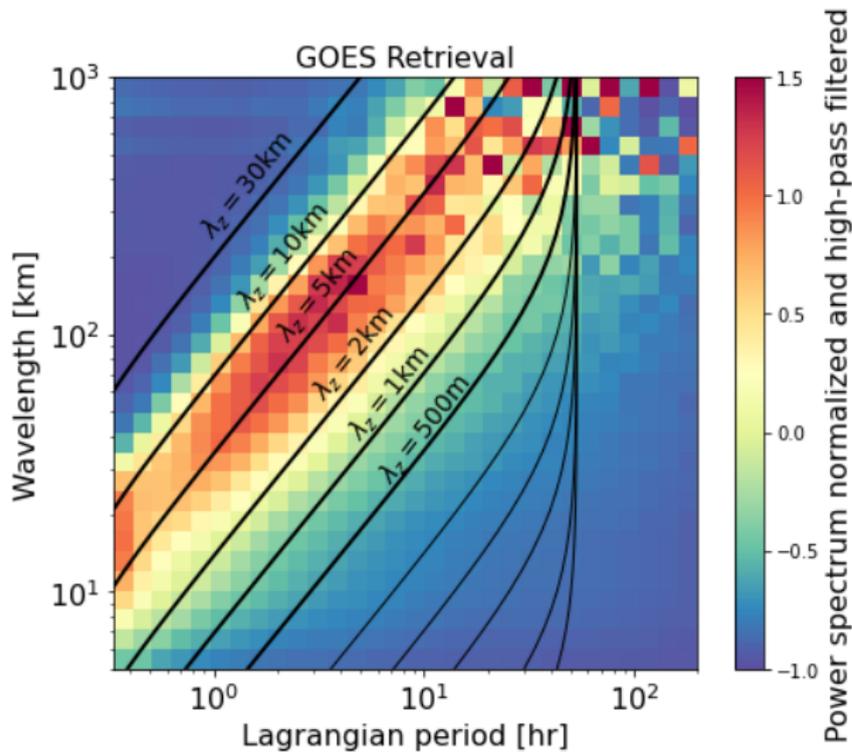


Radiative cooling matters mostly for large scale circulations only in the free troposphere

But local impact of vertical velocity on PBL cooling

## Gravity waves as a good candidate for mesoscale vertical motion

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## Determination of temperature at emission level

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Radiative transfer equation :

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We invert this equation

→ We can deduce temperature  $T^* = f(R_\lambda, T_{sfc})$

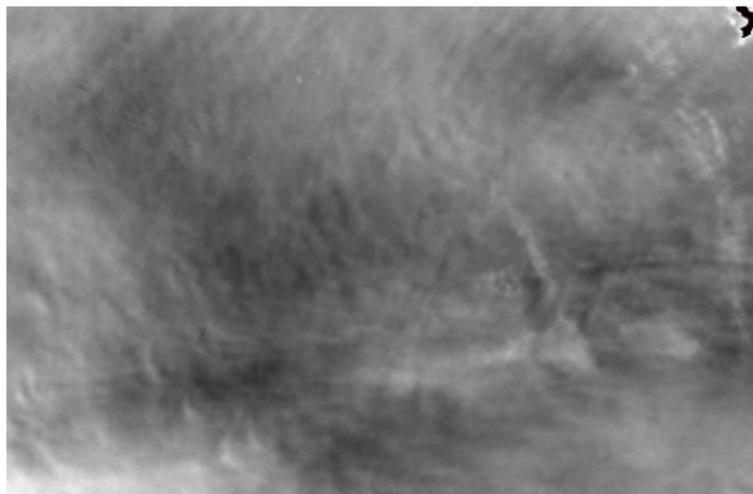
→ Correction from surface or cloud-top temperature inhomogeneities

$$T^* \approx B_\lambda^{-1} \left( \tau^{*\eta} \cdot \frac{R_\lambda - e^{-\tau_s} B_\lambda(T_s)}{\Gamma(1 + \eta) - \tau_s^\eta e^{-\tau_s} \left(1 + \frac{\eta}{\tau_s}\right)} \right)$$

## Determination of horizontal winds

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- High pass filter to select small scale features ( $<30\text{km}$ )
- Band rejection filter at  $10\text{km}$  to remove fast gravity waves



(ex : GOES-E 6.2 $\mu\text{m}$  channel over Barbados, 24 Jan 2020)